



FAKULTA MATEMATIKY,  
FYZIKY A INFORMATIKY

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# Superconducting Parametric Amplifiers

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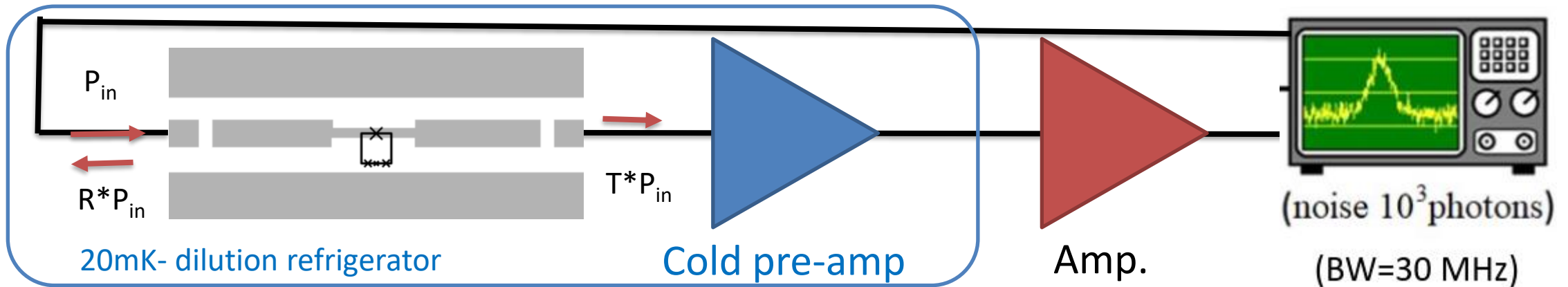


# Experimental requirements for QC

- Circuit QED - qubit state probed by microwave signal - R,T
- Qubit energy  $\Delta/h \sim 6 \text{ GHz} = 300 \text{ mK}$ , for high coherence  $hf \gg k_b T \longrightarrow T < 50 \text{ mK}$
- High-Fidelity Readout  $n_{\text{cav}} \leq 10$  ( $\sim$  expected photons over time  $1/\text{BW}$  in the amp.)

Quantum system, few photon signal

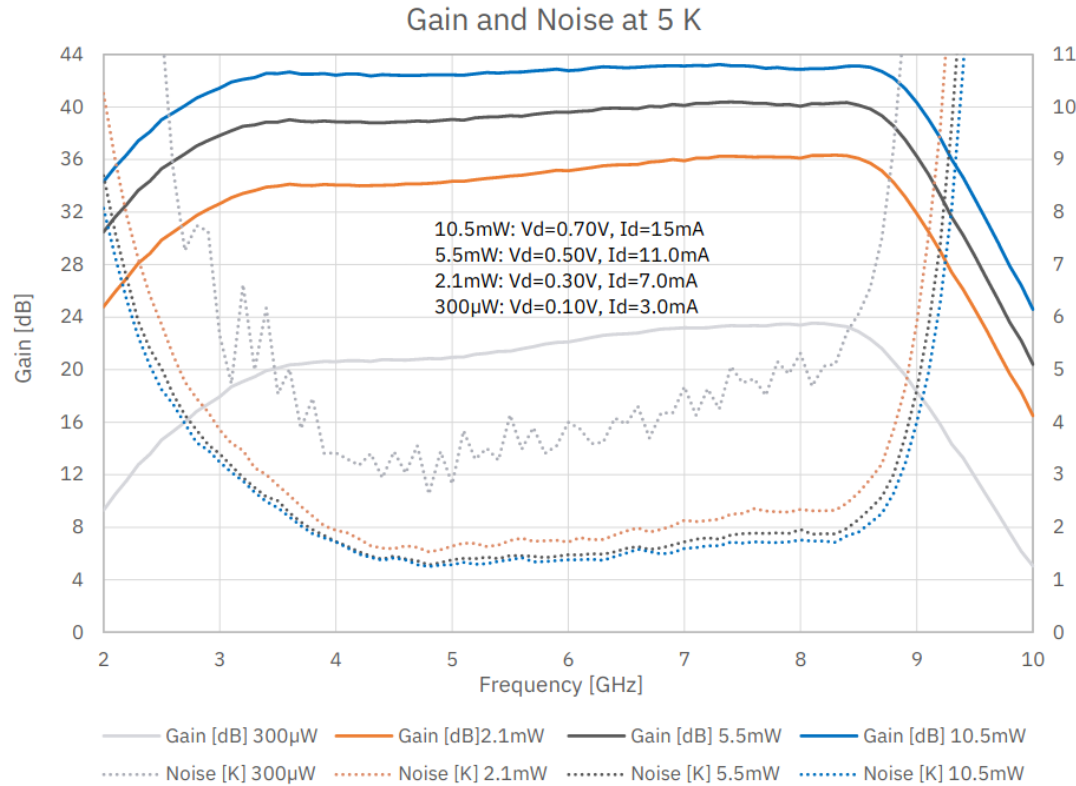
Room temperature electronics



- Amplifier,  $T_N$  equivalent noise temperature - white noise of  $R=50\Omega$ ;  $N_{\text{out}} = G \cdot (N_{\text{in}} + k_b T_N)$
- High signal to noise ratio /low noise temperature is critical! Cryogenic amplifier.

# Cryogenic HEMT amplifiers

- Advantages: Broadband, high gain, flat gain profile, reliable, easy to use, high dynamic range, provide isolation, commercially available



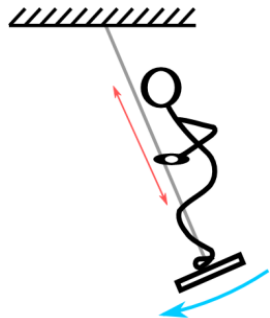
RF Bandwidth	4-8 GHz
Noise Temperature	1.5-2 K
Gain	42 dB
Power dissipation	10 mW

[https://lownoiseactory.com/product/Inf-Inc4\\_8c/](https://lownoiseactory.com/product/Inf-Inc4_8c/)

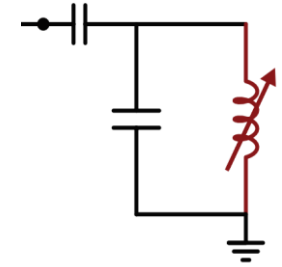
- Disadvantages: Minimum operation temperature 3-4K, high dissipation ~10mW, noise temperature 2 K →  $n_{\text{HEMT}} \sim 10$  photons

# Parametric amplification

- Driven oscillator with a modulated system parameter – energy transfer between modes
  - Child on a swing - resonance frequency  $\omega_0$  modulated (pumped) at  $2\omega_0$

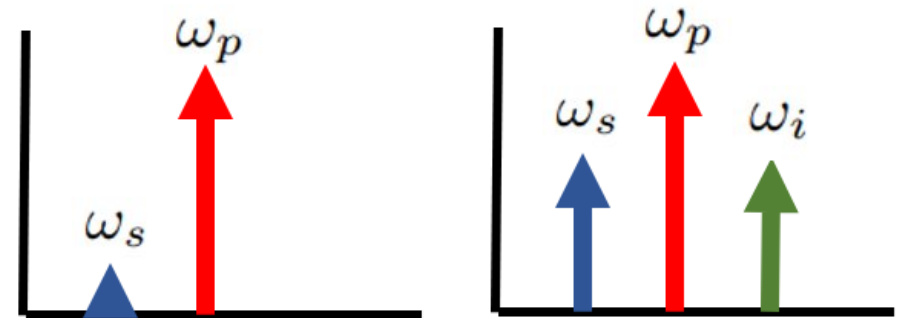
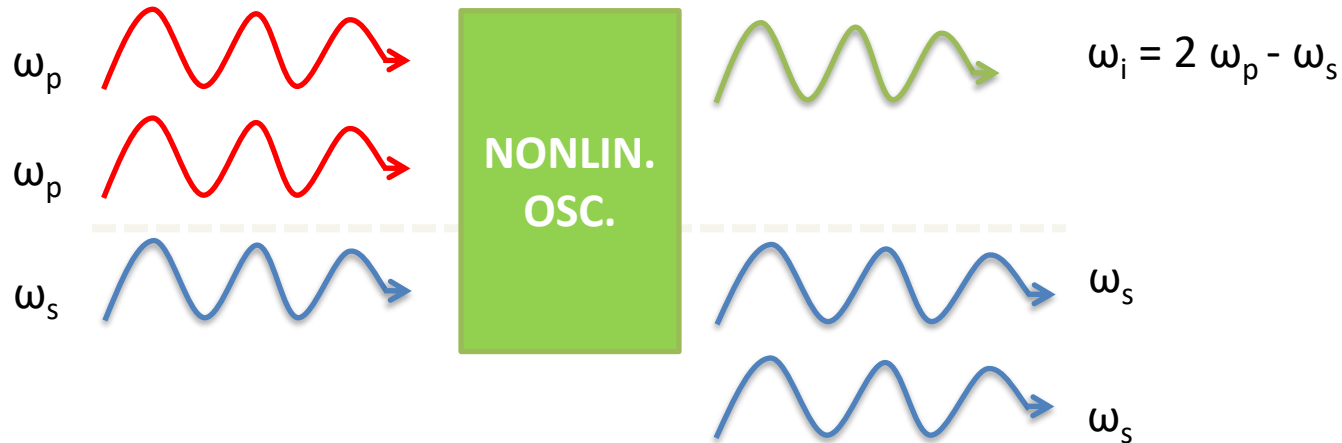


- LC resonator with variable inductance – nonlinear inductance



$$L_0 + \delta L \cos(2\pi f_p t + \phi_p)$$

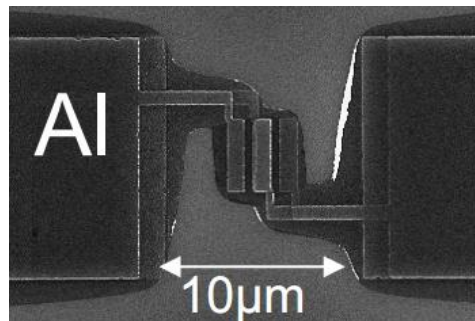
- 4 wave mixing – photon model:  $\omega_p + \omega_p = \omega_s + \omega_i$



# Superconducting parametric amplifiers

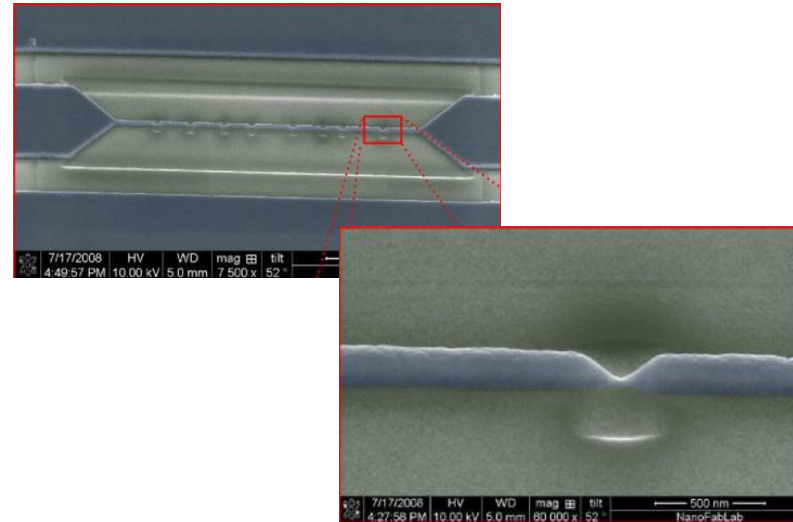
- Nonlinearity – nonlinear inductance

- Josephson Junctions and SQUIDs



$$L(I) = \frac{\Phi_0}{2\pi I_0 \sqrt{1 - (I/I_0)^2}}$$

- Disordered superconductors - High kinetic inductance



E. A. Tholén et al 2009 Phys. Scr. 2009 014019

## Parametric amplification by coupled flux qubits

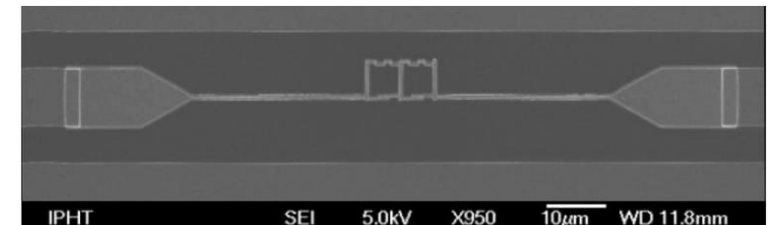
M. Reháč,<sup>1,2</sup> P. Neillinger,<sup>1,2</sup> M. Grajcar,<sup>1,2</sup> G. Oelsner,<sup>3</sup> U. Hübner,<sup>3</sup> E. Il'ichev,<sup>3,4</sup> and H.-G. Meyer<sup>3</sup>

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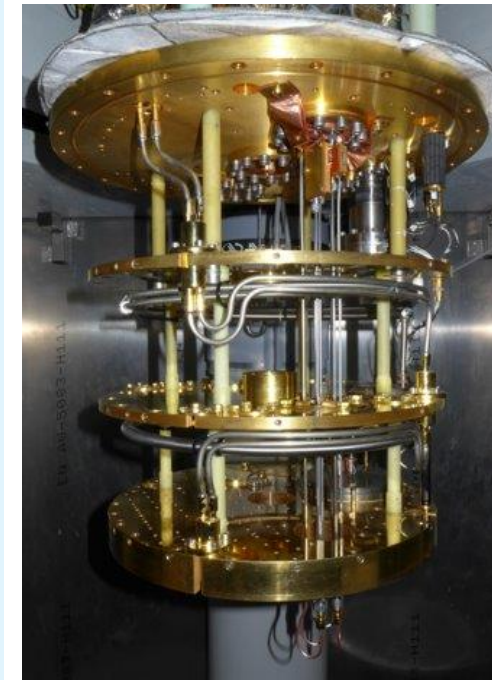
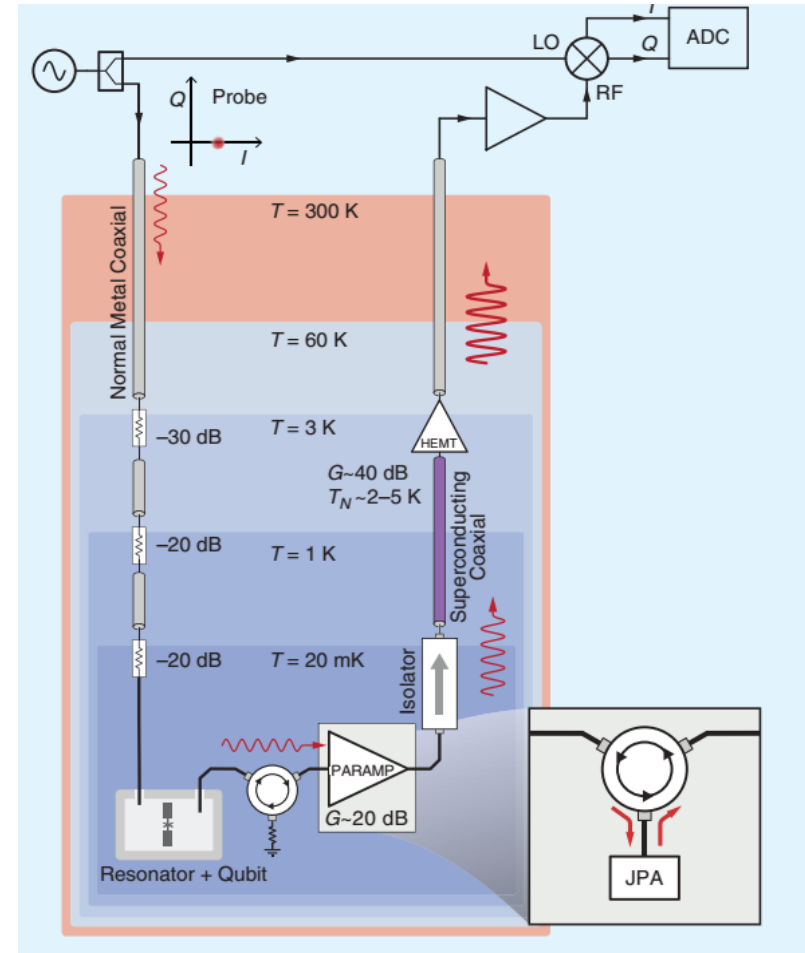
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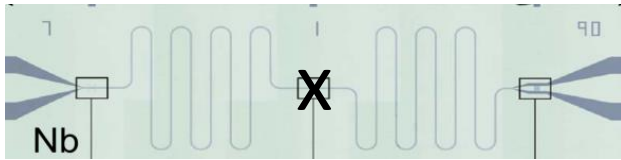
# Experimental set-up for QM with ParAmps

- $^3\text{He}/^4\text{He}$  dilution refrigerator
- Paramps work at 20 mK
- Superconducting ParAmps have no dissipative elements
- Quantum limited amplifier
  - Added-noise number  $n_N^{\text{SQL}} = 1/2$
- Advantages: High gain, ultra-low noise, low power dissipation
- Superconducting ParAmps are an **enabling** technology for superconducting qubit measurement



# Superconducting parametric amplifiers

- Resonant parametric amplifiers

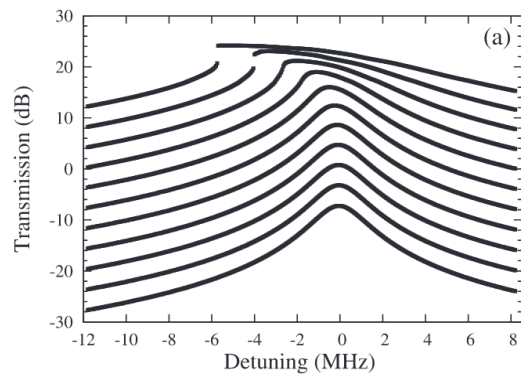


- Duffing oscillator

$$\ddot{y} + \delta\dot{y} + \alpha y + \beta y^3 = F_p \cos(\omega_p t)$$

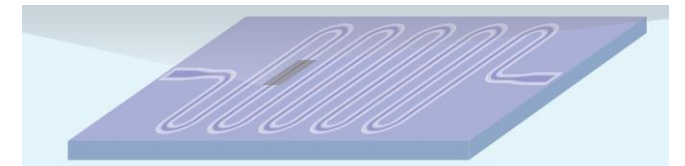
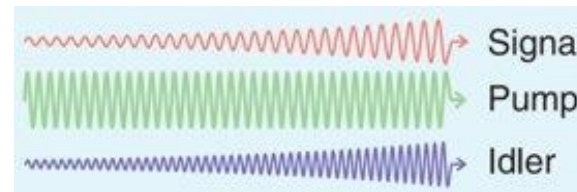
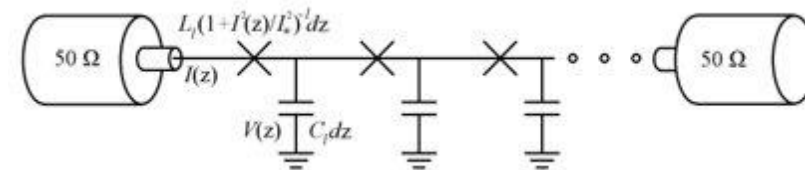
$$y_p = P \cos(\omega_p t - \phi_p)$$

$$P = \frac{F_p}{\sqrt{(\omega_p^2 - \alpha - \frac{3}{4}\beta P^2)^2 + (\delta\omega_p)^2}}$$

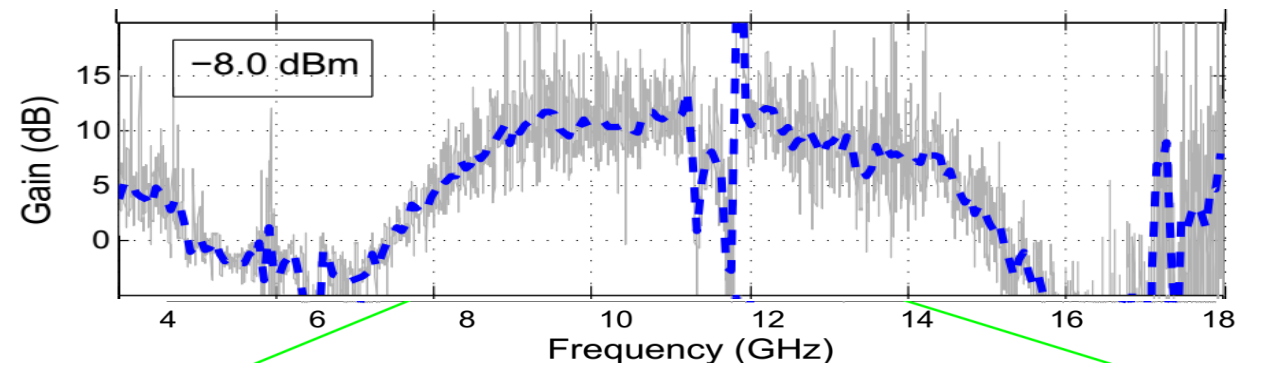


- + SQL noise and below (squeezing), high gain
- Limited bandwidth  $\sim 1$  MHz

- Broadband amplifiers – traveling wave parametric amplifiers TWPAs
- Long nonlinear waveguide, array of JJs  $\rightarrow$  Wave-mixing



J. Aumentado, IEEE Microwave Magazine 21(8):45-59 (2020)



- + Broadband - several GHz
- Gain ripples, demanding fabrication, higher noise

B. H. Eom et al., Nature Physics 8, 623–627 (2012)

# Coupled mode theory

- Current in waveguide

$$I(x, t) = \sum_j^{\{p,s,i\}} \frac{1}{2} \left( I_j(x) e^{i(k_j x - \omega_j t)} + c.c. \right)$$

- Telegrapher's equation for nonlinear medium

$$\frac{\partial^2 I(z, t)}{\partial z^2} - L_l C_l \frac{\partial^2 I(z, t)}{\partial t^2} = \frac{L_l C_l}{6I_c^2} \frac{\partial^2 I(z, t)^3}{\partial t^2}$$

...

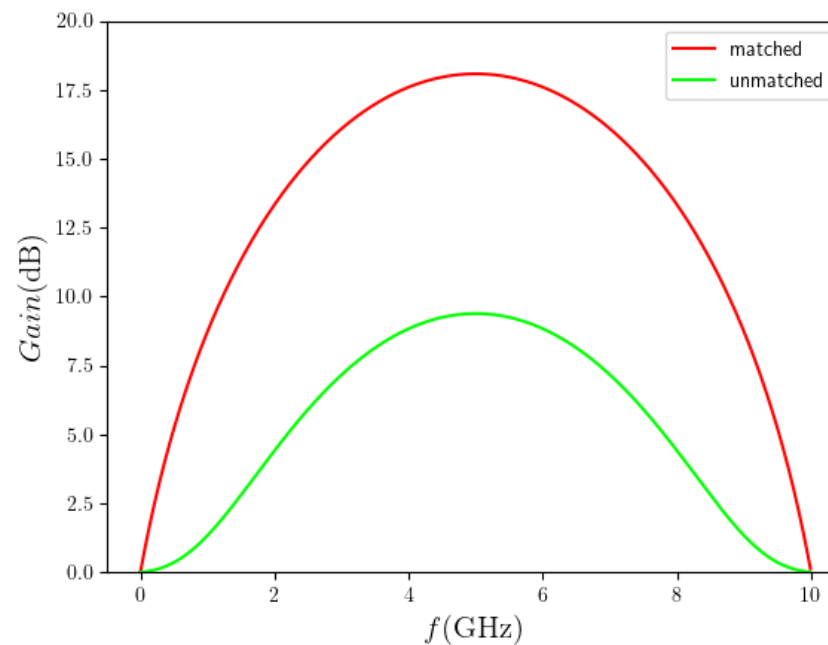
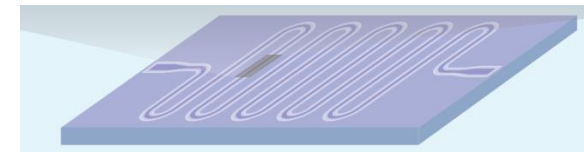
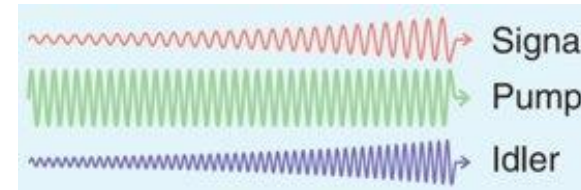
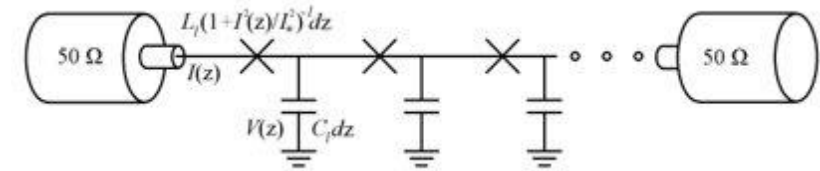
Gain:

$$g = \sqrt{k_s k_i \gamma^2 - \frac{\beta^2}{4}}$$

Phase mismatch:  $\beta = \Delta k(1 + 2\gamma) - 2k_p \gamma$

$$\Delta k = 2k_p - k_s - k_i$$

Nonlinearity strength:  $\gamma = \frac{|t_p I_p|^2}{16I_c^2}$



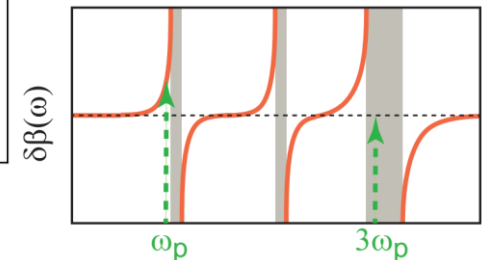
Unmatched phase  $\beta \neq 0$ :

$$G \sim I^2$$

Matched phase  $\beta = 0$ :

$$G \sim \exp(I)$$

Dispersion engineering

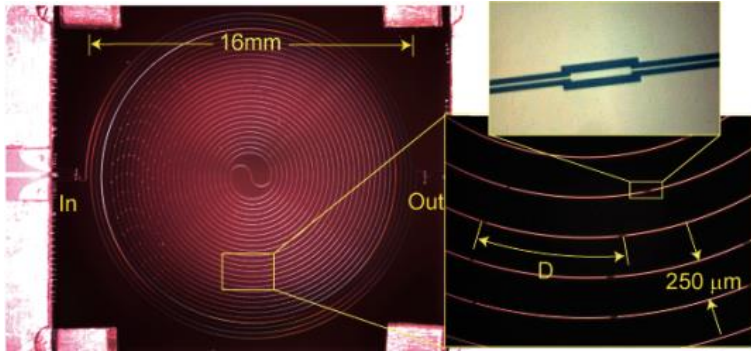


Fixed pump frequency

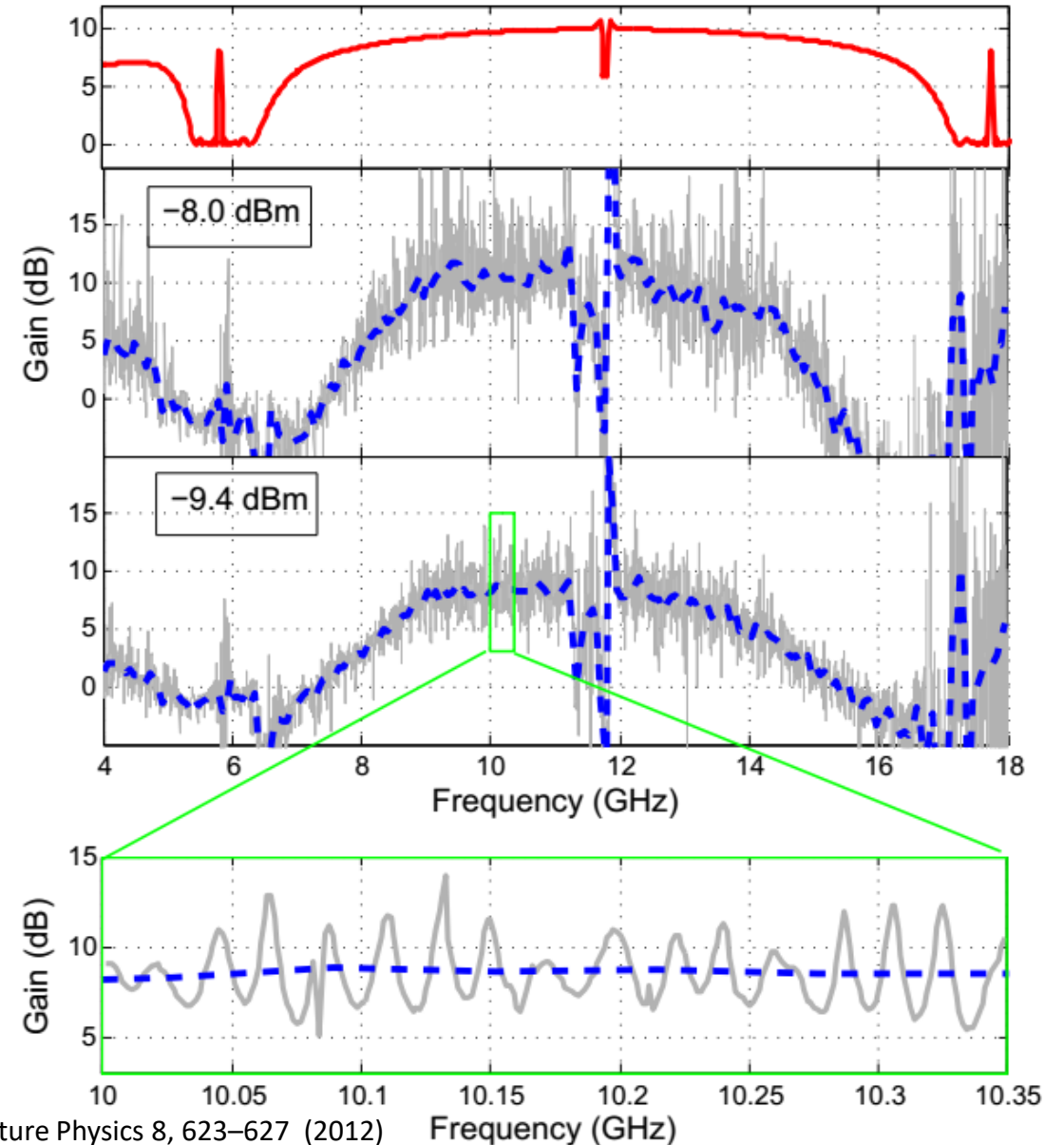


# Coupled mode theory vs experiment

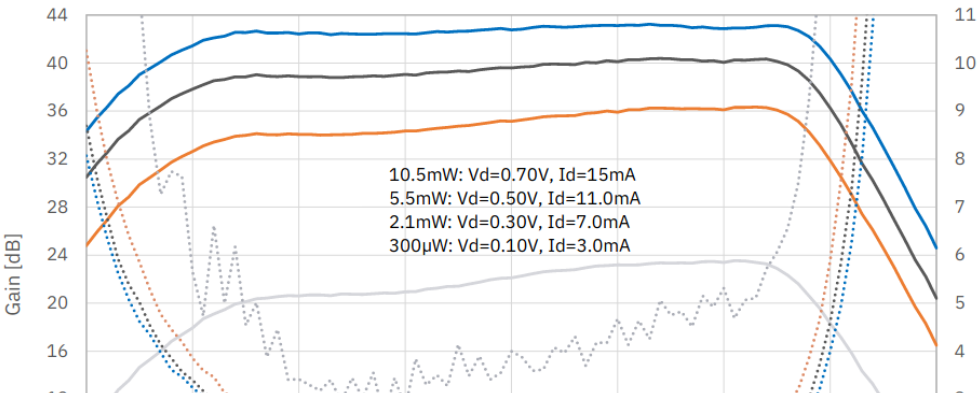
- 80 cm nonlinear waveguide, Impedance  $Z=200 \Omega$



- Phase matched – high gain
- Strong ripples - limits the usability of TWPA
- Coupled mode theory fail to explain the ripples



Gain and Noise at 5 K



# Coupled mode theory with reflections

- Current in finite waveguide with unmatched impedance

$$I(x, t) = \sum_j^{\{p,s,i\}} \frac{1}{2} \left( I_j(x) t_j (e^{ik_j x} + \Gamma e^{-ik_j x}) e^{-i\omega t} + c.c. \right)$$

$$\frac{\partial^2 I(z, t)}{\partial z^2} - L_l C_l \frac{\partial^2 I(z, t)}{\partial t^2} = \frac{L_l C_l}{6I_c^2} \frac{\partial^2 I(z, t)^3}{\partial t^2}$$

Gain: 
$$g = \sqrt{k_s k_i \gamma^2 (1 + 4|\Gamma|^2)} - \frac{\beta}{4}$$

Phase mismatch:

$$\beta = \Delta k (1 + 2\gamma(1 + \Gamma^2)) - 2k_p \gamma (1 - \Gamma^2)$$

Nonlinearity strength: 
$$\gamma = \frac{|t_p I_p|^2}{16I_c^2}$$

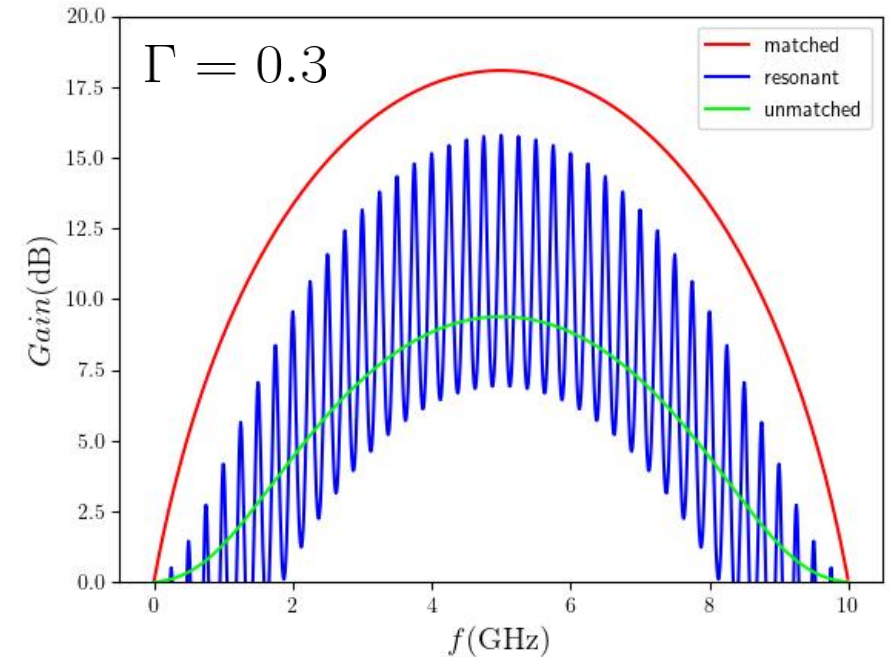
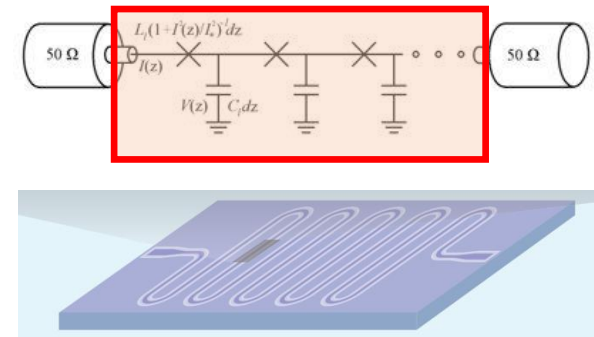
- Fabry-Perot resonator

$Z=200 \Omega$

$$\Gamma = \frac{Z - Z_L}{Z + Z_L}$$

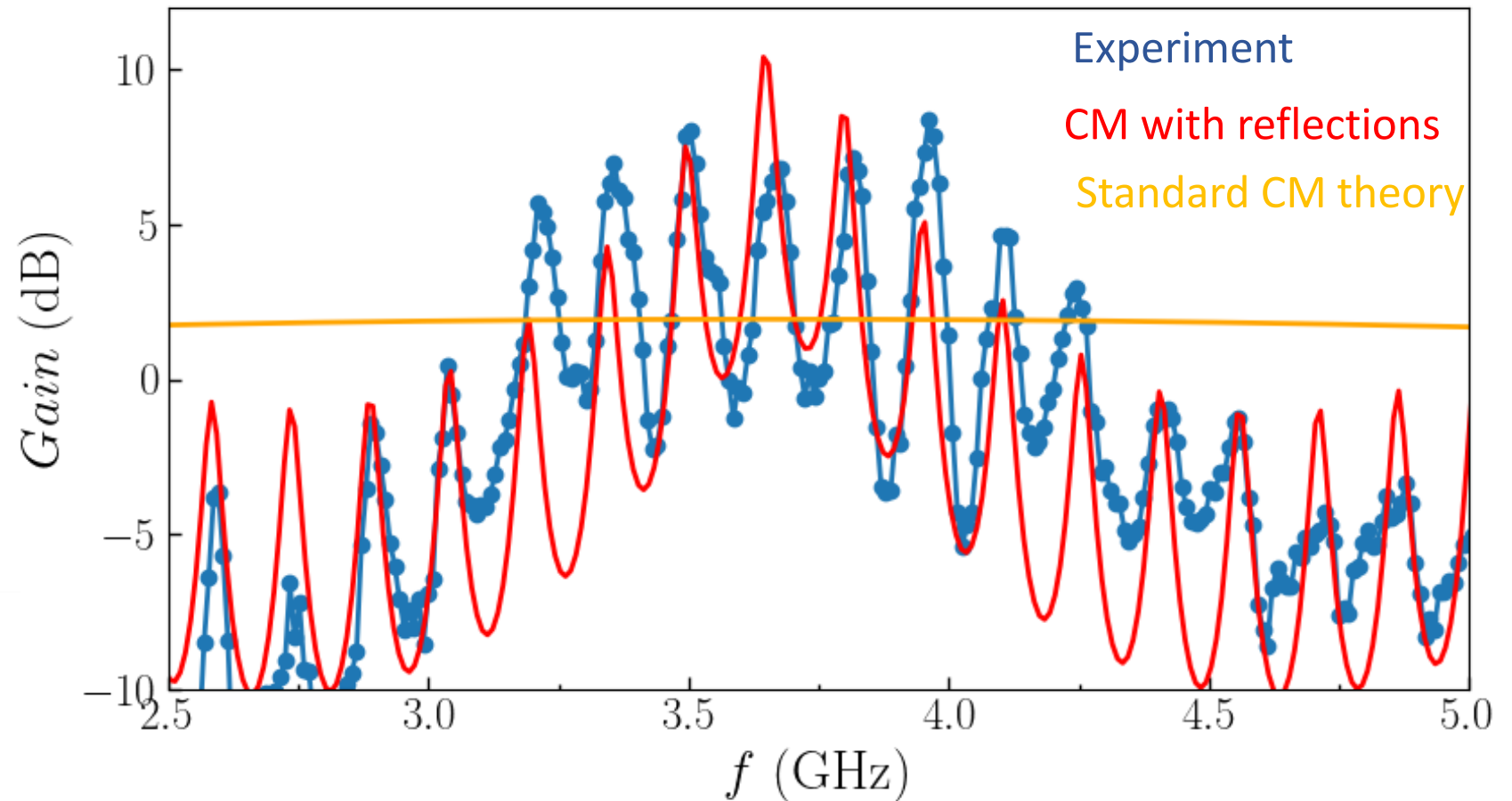
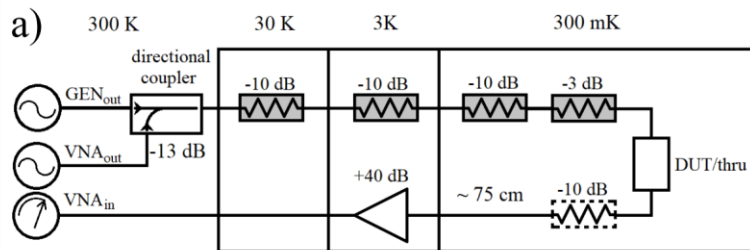
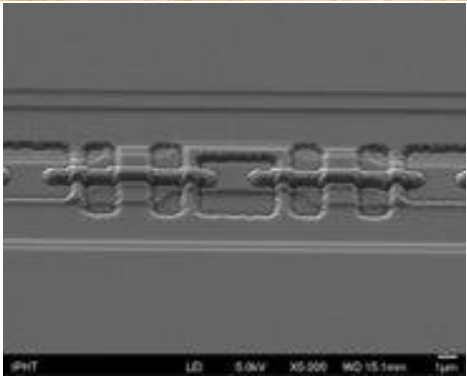
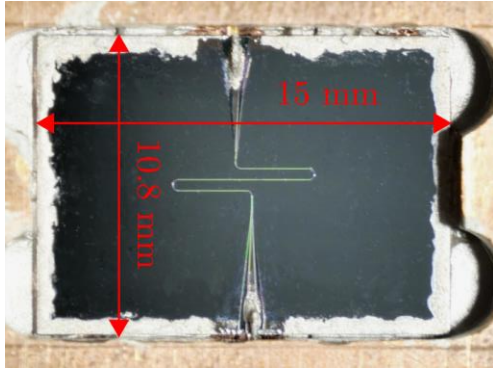
$$t_n = \frac{1}{1 - \Gamma^2 e^{i2k_j L}}$$

$$BW = \frac{\sqrt{1 - \Gamma^2}}{\Gamma} \frac{v}{4\pi l}$$



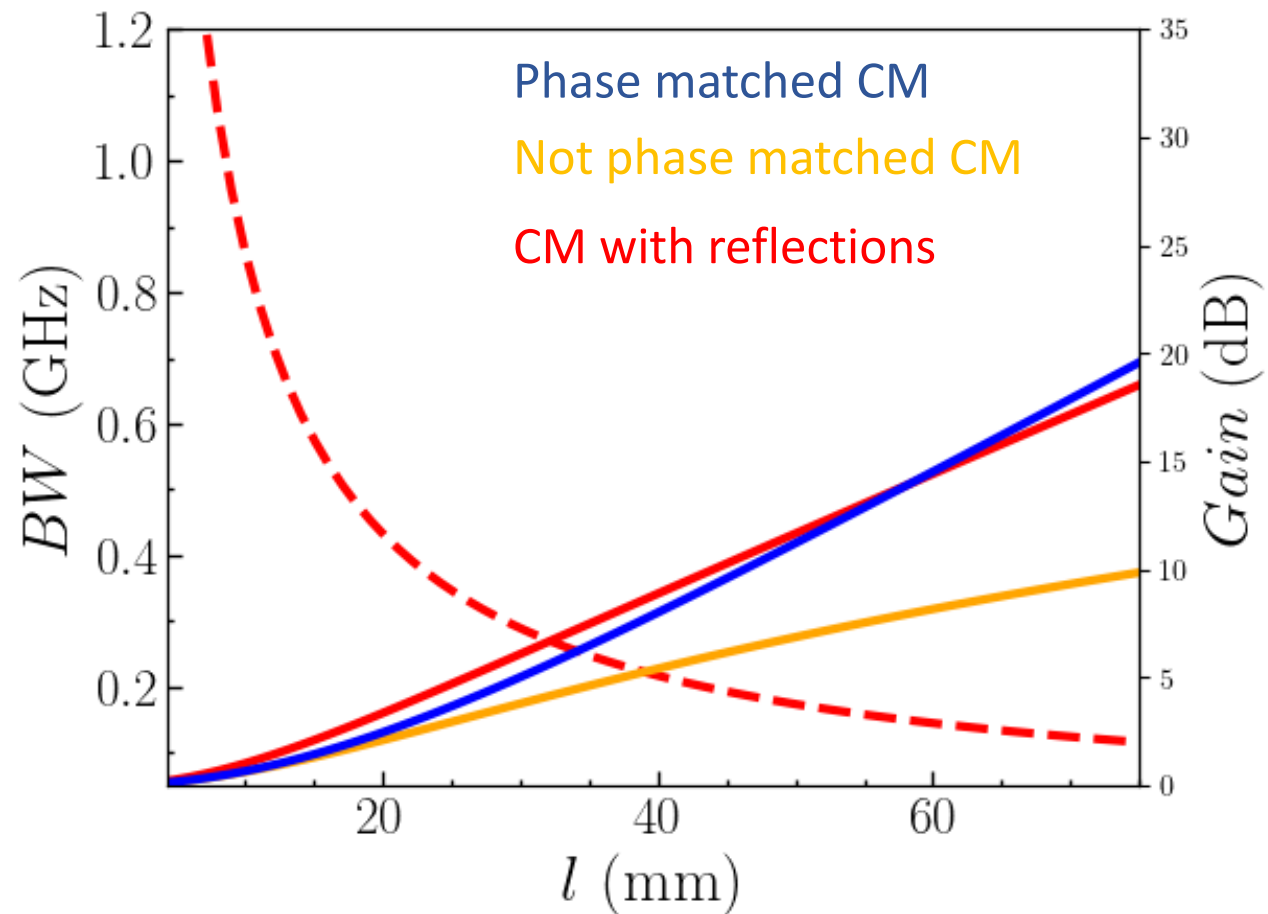
# Nb JJ TWPA - experiment

- Array of 2000 JJ in coplanar waveguide,  $l=11\text{mm}$ ,  $v_{\text{ph}} = 0.14c$ , tested at  $T = 3.5\text{K}$



# Conclusion

- Modified CM theory
  - unmatched TWPA with finite length
- Proper understanding of gain ripples and amplifier bandwidth
- Can be utilized in TWPA design
- The transition regime between traveling-wave and resonant parametric amplifier
- Gain 10dB at 11 mm length without phase matching  
BW~ 100MHz, – tunable design



Thank you for your attention!

S. Kern, et al. arXiv:2203.02448 (2022)